# ASTROFIX BY ALTITUDES OF TWO PAIRS OF STARS ONE NEAR MERIDIAN AND ANOTHER NEAR ELONGATION 

by J. C. Bhattacharji, M.A., F.R.A.S., F.I.G.U. Survey of India


#### Abstract

With a view to improving upon the existing methods of simultaneous determination of latitude and longitude by equal or arbitrary altitudes and of separate determination of latitude by meridian and circum-meridian altitudes of stars, an attempt has been made in this paper to introduce an alternative method of astrofix from observations of only two pairs of stars - one for latitude near meridian transit and the other for longitude (time) near elongation. This method is claimed to be not only simpler and quicker but also more completely free from the usual errors of both vertical collimation and constant of atmospheric refraction.

Apart from its general use in geodetic surveys, the method is also of considerable importance in topographic and hydrographic surveys, and in navigation.


## INTRODUCTION

By astrofix is meant the determination of coordinates of points on the earth's surface by astronomical methods. The precision of astrofix is however dependent on the quality of the angular work and the efficiency of the method of observation actually employed, and is accordingly classified into the following catagories :
(1) First-order astrofix, giving a standard error for the result of $0: 15$ in latitude and $0^{8} .015$ in longitude, obtained from observations on a single night. This is needed for Laplace points in geodetic surveys.
(2) Second-order astrofix, giving a standard error for the result of $1: 0$ in latitude and $0^{\mathrm{s}} .15$ in longitude or less. This is mainly used for astrogeodetic deflections in geodetic surveys but has ample applications in topographic as well as hydrographic surveys.
(3) Third-order astrofix, giving an error for the result of $10^{\prime \prime}$ in latitude and $1^{s}$ in longitude or less. This can be profitably utilised in navigation as well as in explorative surveys.

There are however various methods of determination of astronomical latitudes and longitudes (time) from observations of altitudes of stars. Among these, the more important and now in common use are :
(1) Latitude : The method most generally applied is Talcott, using either a transit or a theodolite provided with a Talcott level. The main drawback of this otherwise theoretically excellent method is that one has often to wait long and to choose stars whose declinations are of only moderate accuracy. As regards the method of circum-meridian altitudes, the computations involved are less simple than in the case of meridian transits. Moreover the results are likely to be affected by the uncertainty in the calculated value of the constant of atmospheric refraction.
(2) Longitude : The usual method is the observation of equal or arbitrary prime vertical altitudes with either a transit or a theodolite fitted with an impersonal micrometer or a Hunter shutter. The main difficulty with the method is that it often keeps the observer waiting indefinitely for completion of the observational programme, due to stars of observation not being always available at all latitudes.
(3) Longitude simultaneous with latitude : The method in common use is the Gauss method. It comprises observations with a prismatic astrolabe of stars of equal altitude. The method of arbitrary altitudes is also gaining ground because of its flexibility. But all the same, the methods have many practical difficulties. As remarked by Roelofs, "A peculiarity of both methods, particularly the Gauss method, is that the observations pass off so easily and smoothly that many observers are inclined to make an exorbitantly large number of observations for safety's sake and realising that they cannot anyhow go to sleep. The computer is thus overwhelmed by this super-abundant material. The obvious remedy is to limit the number of observations to a minimum by a careful pre-selection of stars. It should be added that this does not mean a simple regular distribution of stars in all directions ". Among other disadvantages are long computations, a troublesome programme and loss of accuracy.

An attempt has therefore been made in the present paper to evolve an alternative method of simple and rapid astrofix which can be practically free from the various defects enumerated above, and thus prove of real value to the surveying profession in general.

## LATITUDE

Selection of stars : A star-pair, of magnitudes varying generally from 2.0 to 6.0 depending on the order of accuracy needed, one situated to the north and the other to the south but both within about 15 minutes of their time of meridian transit, must be selected for observation from the fourth fundamental star catalogue FK4. In addition, the altitudes of
the star-pair should preferably be above $30^{\circ}$ and as nearly equal to each other as conveniently possible.

Observational programme : It is always necessary to have a carefully made programme, lasting for about an hour in the case of first-order determinations, about half an hour in the case of second-order determinations and about a quarter of an hour in the case of third-order determinations, showing for the selected star-pair names, aspects, altitudes (computed from the approximate relations : $h=\varphi \pm\left|90^{\circ}-\delta\right|$ for north stars and $h=90^{\circ}-\varphi+\delta$ for south stars), and times (L.S.T.) of passages corresponding to the various azimuth positions of the star-pair (reduced from the relations $a^{\prime \prime}=-15 t^{\beta} \cdot \cos \delta \cdot \sec h$ for north stars and $180^{\circ}-a^{\prime \prime}=$ - $15 t^{s} \cdot \cos \delta \cdot \sec h$ for south stars).

Observational equipment : The astronomical theodolite Wild T4 with accessories, or any other theodolite of similar precision, equipped with a Talcott level, a good chronometer and a stop watch are all that are generally required for the first-order determinations. The Geodetic Tavistock and Wild T3 theodolites are usually used in the second-order determinations. In the case of third-order determinations, a Wild T2 or a Tavistock theodolite should be good enough.

Observational procedure : Before commencing actual observations, the line of horizontal collimation of the theodolite requires to be set in the meridian, correct to a quarter degree or so, knowing only the approximate value of the chronometer error and the chronometer time of meridian transit of a known circumpolar star. In the case where the chronometer error is abnormally large, it can be easily improved by including also a known high altitude meridian star. The routine of observation is that the theodolite is first set at the calculated altitude and azimuth readings of the stars of observation, as per programme, and then the required intersections of stars are made in quick succession with the help of the horizontal wire of the theodolite, noting the corresponding vertical circle readings side by side. In the case of Wild T4 the vertical circle of the theodolite requires to be clamped at the desired vertical settings and then, after carefully noting the respective vertical circle readings, the actual intersections of star-positions are made at each setting in quick succession with the help of the moving micrometer only. The vertical bubble readings are also to be carefully taken immediately after each intersection. The effect of any error in the vertical collimation being completely eliminated in the final results, the observations can be taken with advantage on one face of the theodolite only. For recording instants of intersections use can be made of only the stop watch and the chronometer provided. As regards the series of altitudes measured, these need not be symmetrical with respect to the meridian, but the number of intersections of star-positions for one member, either on the east or the west side of the meridian, should be about the same on the corresponding side of the meridian for the other member of the star-pair.

Computation : Let $h$ be the altitude of a star of declination $\delta$, corresponding to the hour angle $t$ when the star is close to the meridian, and $h_{0}$ the altitude when the star is on the meridian.


Fig. 1
Then on referring to figure 1 we have :
$\sin h=\sin \varphi \cdot \sin \delta+\cos \varphi \cdot \cos \delta \cdot \cos t$
By making use of the Maclaurin formula and following Doolittle's notations, the relation (1) can be reduced to the following forms :
$\varphi=\boldsymbol{h} \mp\left(90^{\circ}-\delta\right)+\mathbf{A} \cdot m-\mathbf{B} \cdot \boldsymbol{n}+\mathbf{C} \cdot o$
for upper (lower) transit of north stars
and
$\varphi=\delta+\left(90^{\circ}-\boldsymbol{h}\right)-\mathbf{A} \cdot \boldsymbol{m}+\mathbf{B} \cdot \boldsymbol{n}-\mathbf{C} \cdot \boldsymbol{o}$
for south stars
where
$\mathrm{A}=\cos \varphi \cdot \cos \delta \cdot \sec h_{0}$
$m=2 \sin ^{2} \frac{t}{2} \cdot \operatorname{cosec} 1^{\prime \prime}$
$\mathrm{B}=\mathrm{A}^{2} \cdot \tan \boldsymbol{h}_{\mathbf{0}}$
$n=2 \sin ^{4} \frac{t}{2} \cdot \operatorname{cosec} 1^{\prime \prime}$
$\mathrm{C}=\mathrm{A}^{3} \cdot \frac{2}{3} \cdot\left(1+3 \tan ^{2} h_{0}\right)$
$o=2 \sin ^{6} \frac{t}{2} \cdot \operatorname{cosec} 1^{\prime \prime}$
or after further simplification,

$$
\begin{align*}
& \varphi=h \mp\left(90^{\circ}-\delta\right)+\mathbf{A} \cdot m-\mathrm{D}+\mathrm{E}  \tag{4}\\
& \text { for upper (lower) transit of north stars }
\end{align*}
$$

and
$\varphi=\delta+\left(90^{\circ}-\boldsymbol{h}\right)-\mathbf{A} \cdot \boldsymbol{m}+\mathbf{D}-\mathbf{E}$
for south stars
where
$\mathbf{A}=\boldsymbol{\operatorname { c o s }} \varphi \cdot \boldsymbol{\operatorname { c o s } \delta \cdot \operatorname { s e c } \boldsymbol { h } _ { \mathbf { 0 } }}$
which is a constant for the particular star of observation;
$m=2 \sin ^{2} \frac{t}{2} \cdot \operatorname{cosec} 1^{\prime \prime}$
which is a tabular quantity directly obtainable from table 1 ;
$\mathrm{D}=2.424 \cdot\left(\frac{\mathrm{~A} \cdot m}{1000}\right)^{2} \cdot \tan h_{0}$
which can be directly read from chart 1 , without requiring actual computations, much more easily and quickly than is normally possible in evaluating $B$ and $n$ separately, and

$$
\mathrm{E}=0.0118\left(\frac{\mathrm{~A} \cdot m}{1000}\right)^{3} \cdot \tan ^{2} h_{0},
$$

which is mostly negligible, but which can be directly read from chart 2 , without recourse to actual computations, much more easily and quickly than is normally possible in evaluating C and $o$ separately.

Since in relations (4) and (5) A is a function of $\varphi$, which is an unknown quantity, and $h_{0}$, which is again a function of $\varphi$, the evaluation of $A$ is not possible without assuming a value $\varphi_{m}$ for $\varphi$, and in that case the computed value $\varphi_{c}$ of $\varphi$, as derived from relation (4) and (5) on substituting $\varphi_{m}$ for $\varphi$ and $\varphi_{m} \pm\left(90^{\circ}-\delta\right)$ for $h_{0}$ in the case of north stars and $\delta+\left(90^{\circ}-\varphi_{m}\right)$ for $h_{0}$ in the case of south stars, will be in error by an amount equal to $\varphi-\varphi_{c}=\Delta \varphi_{c}$ due to :
(1) Error $\Delta \varphi_{m}$ in the assumed value $\varphi_{m}$ of $\varphi$, where $\varphi-\varphi_{m}=\Delta \varphi_{m}$, a constant independent of the star observation.
(2) Error $\Delta h$ in the observed altitudes of stars, on account of the assumed values of both the vertical collimation and the constant of atmospheric refraction. The error is thus practically constant during observations, unless the individual altitudes of the stars of observation differ considerably from each other due to faulty selection of stars.
(3) Error $\Delta t$ in the assumed hour angle $t$, on account of the value of longitude $\lambda$ of the place of observation being an unknown quantity.

But the above errors, when expressed in the form of differential formulae, take the following form :
$\Delta \varphi_{c}=\sec a_{N} \cdot \Delta h-\tan a_{N} \cdot \cos \varphi_{m} \cdot \Delta t$
for north stars
and
$\Delta \varphi_{c}=-\sec a_{\mathrm{S}} \cdot \Delta h+\tan a_{\mathrm{S}} \cdot \cos \varphi_{m} \cdot \Delta t$
for south stars
where $a_{\mathrm{N}}$ and $a_{\mathrm{S}}$ denote small azimuthal angles as approximately derived from the relations :
$a_{N}^{\prime \prime}=-15 \cdot t^{s} \cdot \cos \delta \cdot \sec h_{0}$ and $180^{\circ}-a_{S}^{\prime \prime}=-15 \cdot t^{s} \cdot \cos \delta \cdot \sec h_{0}$ and suffixes $N$ and $S$ refer to north and south stars respectively.

TABLE I
for
values of $m: 2 \sin ^{2} \frac{t}{2} \cdot \operatorname{cosec} 1^{\prime \prime}$

| Hour angle t | $0^{\text {m }}$ | $1^{\text {m }}$ | $2^{m}$ | $3^{m}$ | $4^{m}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $0^{\text {s }}$ | 0! 00 | 1!96 | 7. 85 | $17!67$ | 31.42 |
| 1 | 0.00 | 2.03 | 7.98 | 17.87 | 31.68 |
| 2 | 0.00 | 2. 10 | 8.12 | 18.07 | 31.94 |
| 3 | 0.00 | 2.16 | 8.25 | 18.27 | 32.20 |
| 4 | 0.01 | 2.23 | 8.39 | 18.47 | 32.47 |
| 5 | 0.01 | 2.31 | 8.52 | 18.67 | 32.74 |
| 6 | 0.02 | 2.38 | 8.66 | 18.87 | 33.01 |
| 7 | 0. 02 | 2.45 | 8.80 | 19.07 | 33.27 |
| 8 | 0. 03 | 2.52 | 8. 94 | 19.28 | 33.54 |
| 9 | 0. 04 | 2. 60 | 9.08 | 19.48 | 33.81 |
| 10 | 0.05 | 2.67 | 9.22 | 19.69 | 34.09 |
| 11 | 0.06 | 2. 75 | 9.36 | 19.90 | 34.36 |
| 12 | 0.08 | 2.83 | 9.50 | 20.11 | 34.64 |
| 13 | 0.09 | 2.91 | 9.64 | 20.32 | 34.91 |
| 14 | 0.11 | 2.99 | 9.79 | 20.63 | 35.19 |
| 15 | 0.12 | 3. 07 | 9.94 | 20. 74 | 35.46 |
| 16 | 0.14 | 3.15 | 10.09 | 20.95 | 35.74 |
| 17 | 0.16 | 3.23 | 10.24 | 21.16 | 36.02 |
| 18 | 0. 18 | 3.32 | 10.39 | 21.38 | 36.30 |
| 19 | 0. 20 | 3. 40 | 10.54 | 21.60 | 36.58 |
| 20 | 0. 22 | 3. 49 | 10.69 | 21.82 | 36.87 |
| 21 | 0. 24 | 3.58 | 10.84 | 22.03 | 37.15 |
| 22 | 0.26 | 3. 67 | 11.00 | 22.25 | 37.44 |
| 23 | 0.28 | 3.76 | 11.15 | 22.47 | 37. 72 |
| 24 | 0.31 | 3. 85 | 11.31 | 22.70 | 38. 01 |
| 25 | 0. 34 | 3.94 | 11.47 | 22.92 | 38.30 |
| 26 | 0.37 | 4.03 | 11.63 | 23.14 | 38.59 |
| 27 | 0.40 | 4.12 | 11.79 | 23.37 | 38.88 |
| 28 | 0.43 | 4.22 | 11.95 | 23.60 | 39.17 |
| 29 | 0.46 | 4.32 | 12.11 | 23.82 | 39.46 |
| 30 | 0.49 | 4.42 | 12.27 | 24.05 | 39.76 |
| 31 | 0.52 | 4.52 | 12.43 | 24.28 | 40.05 |
| 32 | 0.56 | 4.62 | 12.60 | 24.51 | 40.35 |
| 33 | 0.59 | 4.72 | 12.76 | 24.74 | 40.65 |
| 34 | 0.63 | 4.82 | 12.93 | 24.98 | 40.95 |
| 35 | 0.67 | 4.92 | 13.10 | 25.21 | 41.25 |
| 36 | 0. 71 | 5.03 | 13.27 | 25.45 | 41.55 |
| 37 | 0. 75 | 5.13 | 13.44 | 25.68 | 41.85 |
| 38 | 0. 79 | 5.24 | 13.62 | 25.92 | 42.15 |
| 39 | 0. 83 | 5.34 | 13.79 | 26.16 | 42.45 |
| 40 | 0. 87 | 5.45 | 13.96 | 26.40 | 42.76 |
| 41 | 0. 91 | 5.56 | 14.13 | 26.64 | 43.06 |
| 42 | 0.96 | 5.67 | 14.31 | 26.88 | 43.37 |
| 43 | 1. 01 | 5. 78 | 14.49 | 27.12 | 43.68 |
| 44 | 1.06 | 5.90 | 14.67 | 27.37 | 43.99 |


| Hour angle <br> $\mathbf{t}$ | $\mathbf{0}^{\mathrm{m}}$ | $1^{\mathrm{m}}$ | $2^{m}$ | $3^{m}$ | $4^{m}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |
| 45 | 1.10 | 6.01 | 14.85 | 27.61 | 44.30 |
| 46 | 1.15 | 6.13 | 15.03 | 27.86 | 44.61 |
| 47 | 1.20 | 6.24 | 15.21 | 28.10 | 44.92 |
| 48 | 1.26 | 6.36 | 15.39 | 28.35 | 45.24 |
| 49 | 1.31 | 6.48 | 15.57 | 28.60 | 45.55 |
| 50 | 1.36 | 6.60 | 15.76 | 28.85 | 45.87 |
| 51 | 1.42 | 6.72 | 15.95 | 29.10 | 46.18 |
| 52 | 1.48 | 6.84 | 16.14 | 29.36 | 46.50 |
| 53 | 1.53 | 6.96 | 16.32 | 29.61 | 46.82 |
| 54 | 1.59 | 7.09 | 16.51 | 29.86 | 47.14 |
| 55 | 1.65 | 7.21 | 16.70 | 30.12 | 47.46 |
| 56 | 1.71 | 7.34 | 16.89 | 30.38 | 47.79 |
| 57 | 1.77 | 7.46 | 17.08 | 30.64 | 48.11 |
| 58 | 1.83 | 7.60 | 17.28 | 30.90 | 48.43 |
| 59 | 1.89 | 7.72 | 17.47 | 31.16 | 48.76 |
| 60 | 1.96 | 7.85 | 17.67 | 31.42 | 49.09 |
|  |  |  |  |  |  |


| Oiff. per <br> $1^{s}$ | 0.1 | $0^{s} .2$ | 0.3 | 0.4 | $0^{s} .5$ | 0.6 | 0.7 | 0.8 | 0.9 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.02 | 0.00 | 0.00 | 0.01 | 0.01 | 0.01 | 0.401 | 0.01 | 0.02 | 0.02 |
| 0.04 | 0.00 | 0.01 | 0.01 | 0.02 | 0.02 | 0.02 | 0.03 | 0.03 | 0.04 |
| 0.06 | 0.01 | 0.01 | 0.02 | 0.02 | 0.03 | 0.04 | 0.04 | 0.05 | 0.05 |
| 0.08 | 0.01 | 0.02 | 0.02 | 0.03 | 0.04 | 0.05 | 0.06 | 0.06 | 0.07 |
| 0.10 | 0.01 | 0.02 | 0.03 | 0.04 | 0.05 | 0.06 | 0.07 | 0.08 | 0.09 |
| 0.12 | 0.01 | 0.02 | 0.04 | 0.05 | 0.06 | 0.07 | 0.08 | 0.10 | 0.11 |
| 0.14 | 0.01 | 0.03 | 0.04 | 0.06 | 0.07 | 0.08 | 0.10 | 0.11 | 0.13 |
| 0.16 | 0.02 | 0.03 | 0.05 | 0.06 | 0.08 | 0.10 | 0.11 | 0.13 | 0.14 |
| 0.18 | 0.02 | 0.04 | 0.05 | 0.07 | 0.09 | 0.11 | 0.13 | 0.14 | 0.16 |
| 0.20 | 0.02 | 0.04 | 0.06 | 0.08 | 0.10 | 0.12 | 0.14 | 0.16 | 0.18 |
| 0.22 | 0.02 | 0.04 | 0.07 | 0.09 | 0.11 | 0.13 | 0.15 | 0.18 | 0.20 |
| 0.24 | 0.02 | 0.05 | 0.07 | 0.10 | 0.12 | 0.14 | 0.17 | 0.19 | 0.22 |
| 0.26 | 0.03 | 0.05 | 0.08 | 0.10 | 0.13 | 0.16 | 0.18 | 0.21 | 0.23 |
| 0.28 | 0.03 | 0.06 | 0.08 | 0.11 | 0.14 | 0.17 | 0.20 | 0.22 | 0.25 |
| 0.30 | 0.03 | 0.06 | 0.09 | 0.12 | 0.15 | 0.18 | 0.21 | 0.24 | 0.27 |
| 0.32 | 0.03 | 0.06 | 0.10 | 0.13 | 0.16 | 0.19 | 0.22 | 0.26 | 0.29 |


| Hour angle | $5^{\text {m }}$ | $6^{\text {m }}$ | $7^{m}$ | $8^{n}$ | $9^{\text {m }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $0^{5}$ | 49:09 | 70.'68 | 96!20 | 125!'65 | 159.'0 |
| 1 | 49.41 | 71.07 | 96.66 | 126.17 | 159.6 |
| 2 | 49.74 | 71.47 | 97.12 | 126.70 | 160.2 |
| 3 | 50.07 | 71.86 | 97.58 | 127. 22 | 160.8 |
| 4 | 50.40 | 72.26 | 98. 04 | 127.75 | 161.3 |
| 5 | 50.73 | 72.66 | 98.50 | 128.28 | 161.9 |
| 6 | 51.07 | 73.06 | 98.97 | 128.81 | 162.5 |
| 7 | 51.40 | 73.46 | 99.43 | 129.34 | 163.1 |
| 8 | 51.74 | 73.86 | 99.90 | 129.87 | 163.7 |
| 9 | 52.07 | 74.26 | 100.37 | 130.40 | 164.3 |
| 10 | 52.41 | 74.66 | 100.84 | 130.94 | 164.9 |
| 11 | 52.75 | 75.06 | 101.31 | 131.47 | 165.5 |
| 12 | 53.09 | 75.47 | 101.78 | 132.01 | 166.1 |
| 13 | 53.43 | 75.88 | 102.25 | 132.55 | 166.7 |
| 14 | 53.77 | 76.29 | 102. 72 | 133.09 | 167.3 |
| 15 | 54.11 | 76.69 | 103.20 | 133.63 | 167.9 |
| 16 | 54.46 | 77.10 | 103.67 | 134.17 | 168.5 |
| 17 | 54.80 | 77.51 | 104.15 | 134.71 | 169.1 |
| 18 | 55.15 | 77.93 | 104.63 | 135.25 | 169.8 |
| 19 | 55.50 | 78.34 | 105.10 | 135.80 | 170.4 |
| 20 | 55.84 | 78.75 | 105.58 | 136.34 | 171.0 |
| 21 | 56.19 | 79.16 | 106.06 | 136.88 | 171.6 |
| 22 | 56.55 | 79.58 | 106.55 | 137.43 | 172.2 |
| 23 | 56.90 | 80.00 | 107.03 | 137.98 | 172.85 |
| 24 | 57.25 | 80.42 | 107.51 | 138.53 | 173.4 |
| 25 | 57.60 | 80.84 | 107.99 | 139.08 | 174.0 |
| 26 | 57.96 | 81.26 | 108.48 | 139.63 | 174.70 |
| 27 | 58.32 | 81.68 | 108.97 | 140.18 | 175.32 |
| 28 | 58.68 | 82.10 | 109.46 | 140.74 | 175.9 |
| 29 | 59.03 | 82.52 | 109.95 | 141.29 | 176.5 |
| 30 | 59.40 | 82.95 | 110.44 | 141.85 | 177.18 |
| 31 | 59.75 | 83.38 | 110.93 | 142.40 | 177.80 |
| 32 | 60.11 | 83.81 | 111.43 | 142.96 | 178.4 |
| 33 | 60.47 | 84.23 | 111.92 | 143.52 | 179.05 |
| 34 | 60.84 | 84.66 | 112.41 | 144.08 | 179.68 |
| 35 | 61.20 | 85.09 | 112.90 | 144.64 | 180.30 |
| 36 | 61.57 | 85.52 | 113.40 | 145.20 | 180.93 |
| 37 | 61.94 | 85.95 | 113.90 | 145.76 | 181.56 |
| 38 | 62.31 | 86.39 | 114.40 | 146.33 | 182.19 |
| 39 | 62.68 | 86.82 | 114.90 | 146.89 | 182.82 |
| 40 | 63.05 | 87.26 | 115.40 | 147.46 | 183.46 |
| 41 | 63.42 | 87.70 | 115.90 | 148.03 | 184.09 |
| 42 | 63.79 | 88.14 | 116.40 | 148.60 | 184.7 |
| 43 | 64.16 | 88.57 | 116.90 | 149.17 | 185.35 |
| 44 | 64.54 | 89.01 | 117.41 | 149.74 | 185.99 |
| 45 | 64.91 | 89.45 | 117.92 | 150.31 | 186.63 |
| 46 | 65.29 | 89.89 | 118.43 | 150.88 | 187.27 |
| 47 | 65.67 | 90.33 | 118.94 | 151.45 | 187.91 |
| 48 | 66.05 | 90. 78 | 119.45 | 152.03 | 188.55 |


| Iour angle | $5^{m}$ | $6^{m}$ | $7^{m}$ | $8^{m}$ | $9^{m}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $t$ |  | 66.43 | 91.23 | 119.96 | 152.61 |
| 49 | 66.81 | 91.68 | 120.47 | 153.19 | 189.19 |
| 50 | 67.19 | 92.12 | 120.98 | 153.77 | 189.83 |
| 51 | 67.58 | 92.57 | 121.49 | 154.35 | 190.47 |
| 52 | 67.96 | 93.02 | 122.01 | 154.93 | 191.76 |
| 53 | 68.35 | 93.47 | 122.53 | 155.51 | 192.41 |
| 54 | 68.73 | 93.92 | 123.05 | 156.09 | 193.06 |
| 55 | 69.12 | 94.38 | 123.57 | 156.67 | 193.71 |
| 56 | 69.51 | 94.83 | 124.09 | 157.25 | 194.36 |
| 57 | 69.90 | 95.29 | 124.61 | 157.84 | 195.01 |
| 58 | 70.29 | 95.74 | 125.13 | 158.43 | 195.66 |
| 59 | 70.68 | 96.20 | 125.65 | 159.02 | 196.32 |
| 60 |  |  |  |  |  |


| Diff. per <br> $1^{s}$ | $0 . s 1$ | 0.2 | 0.3 | 0.4 | 0.5 | 0.6 | 0.7 | 0.8 | 0.9 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $0!32$ | 0.03 | 0.06 | 0.110 | 0.113 | 0.16 | 0.19 | 0.22 | 0.26 | 0.29 |
| 0.34 | 0.03 | 0.07 | 0.10 | 0.14 | 0.17 | 0.20 | 0.24 | 0.27 | 0.31 |
| 0.36 | 0.04 | 0.07 | 0.11 | 0.14 | 0.18 | 0.22 | 0.25 | 0.29 | 0.32 |
| 0.38 | 0.04 | 0.08 | 0.11 | 0.15 | 0.19 | 0.23 | 0.27 | 0.30 | 0.34 |
| 0.40 | 0.04 | 0.08 | 0.12 | 0.16 | 0.20 | 0.24 | 0.28 | 0.32 | 0.36 |
| 0.42 | 0.04 | 0.08 | 0.13 | 0.17 | 0.21 | 0.25 | 0.29 | 0.34 | 0.38 |
| 0.44 | 0.04 | 0.09 | 0.13 | 0.18 | 0.22 | 0.26 | 0.31 | 0.35 | 0.40 |
| 0.46 | 0.05 | 0.09 | 0.14 | 0.18 | 0.23 | 0.28 | 0.32 | 0.37 | 0.41 |
| 0.48 | 0.05 | 0.10 | 0.14 | 0.19 | 0.24 | 0.29 | 0.34 | 0.38 | 0.43 |
| 0.50 | 0.05 | 0.10 | 0.15 | 0.20 | 0.25 | 0.30 | 0.35 | 0.40 | 0.45 |
| 0.52 | 0.05 | 0.10 | 0.16 | 0.21 | 0.26 | 0.31 | 0.36 | 0.42 | 0.47 |
| 0.54 | 0.05 | 0.11 | 0.16 | 0.22 | 0.27 | 0.32 | 0.38 | 0.43 | 0.49 |
| 0.56 | 0.06 | 0.11 | 0.17 | 0.22 | 0.28 | 0.34 | 0.39 | 0.45 | 0.50 |
| 0.58 | 0.06 | 0.12 | 0.17 | 0.23 | 0.29 | 0.35 | 0.41 | 0.46 | 0.52 |
| 0.60 | 0.06 | 0.12 | 0.18 | 0.24 | 0.30 | 0.36 | 0.42 | 0.48 | 0.54 |
| 0.62 | 0.06 | 0.12 | 0.19 | 0.25 | 0.31 | 0.37 | 0.43 | 0.50 | 0.56 |
| 0.64 | 0.06 | 0.13 | 0.19 | 0.26 | 0.32 | 0.38 | 0.45 | 0.51 | 0.58 |


| Hour angle | $10^{m}$ | $11^{\text {m }}$ | $12^{m}$ | $13^{m}$ | $14^{\text {m }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $0^{5}$ | 196.'32 | 237.54 | 282.' 68 | 331". 74 | 384". 7 |
| 1 | 196.97 | 238.26 | 283.47 | 332.59 | 385.6 |
| 2 | 197.63 | 238.98 | 284.26 | 333.44 | 386.5 |
| 3 | 198.28 | 239.70 | 285.04 | 334.29 | 387.4 |
| 4 | 198.94 | 240.42 | 285.83 | 335.15 | 388.4 |
| 5 | 199.60 | 241.14 | 286.62 | 336.00 | 389. 3 |
| 6 | 200.26 | 241.87 | 287. 41 | 336.86 | 390.2 |
| 7 | 200.92 | 242.60 | 288.20 | 337. 72 | 391.1 |
| 8 | 201.59 | 243.33 | 289.00 | 338.58 | 392.09 |
| 9 | 202.25 | 244.06 | 289.79 | 339.44 | 393.0 |
| 10 | 202.32 | 244.79 | 230.58 | 340.30 | 333.3 |
| 11 | 203.58 | 245.52 | 291.38 | 341.16 | 394.8 |
| 12 | 204.25 | 246.25 | 292.18 | 342.02 | 395.7 |
| 13 | 204.92 | 246.98 | 292.98 | 342.88 | 396.7 |
| 14 | 205.59 | 247.72 | 293. 78 | 343.75 | 397.65 |
| 15 | 206.26 | 248.45 | 294.58 | 344.62 | 398.5 |
| 16 | 206.93 | 249.19 | 295.38 | 345.49 | 399.5 |
| 17 | 207.60 | 249.93 | 296. 18 | 346.36 | 400.45 |
| 18 | 208.27 | 250.67 | 296.99 | 347.23 | 401.3 |
| 19 | 208. 94 | 251.41 | 297.79 | 348.10 | 402.32 |
| 20 | 209.62 | 252.15 | 298. 60 | 348.97 | 403.26 |
| 21 | 210.30 | 252.89 | 299.40 | 349.84 | 404.20 |
| 22 | 210.98 | 253.63 | 300.21 | 350.71 | 405.14 |
| 23 | 211.66 | 254.37 | 301.02 | 351.58 | 406.08 |
| 24 | 212.34 | 255.12 | 301.83 | 352.46 | 407.02 |
| 25 | 213.02 | 255.87 | 302.64 | 353.34 | 407.9 |
| 26 | 213.70 | 256.62 | 303.46 | 354.22 | 408.90 |
| 27 | 214.38 | 257.37 | 304.27 | 355.10 | 409.8 |
| 28 | 215.07 | 258.12 | 305.09 | 355.98 | 410.73 |
| 29 | 215.75 | 258.87 | 305.90 | 356.86 | 411.73 |
| 30 | 216.44 | 259.62 | 306.72 | 357.74 | 412.68 |
| 31 | 217.12 | 260.37 | 307.54 | 358.62 | 413.63 |
| 32 | 217.81 | 261.12 | 308.36 | 359.51 | 414.59 |
| 33 | 218.50 | 261.88 | 309.18 | 360.39 | 415.54 |
| 34 | 219.19 | 262.64 | 310.00 | 361.28 | 416.49 |
| 35 | 219.88 | 263.39 | 310.82 | 362.17 | 417.44 |
| 36 | 220.58 | 264.15 | 311.65 | 363. 07 | 418.40 |
| 37 | 221.27 | 26.4. 91 | 312.47 | 363.96 | 419.35 |
| 38 | 221.97 | 265.68 | 313.30 | 364.85 | 420.31 |
| 39 | 222.66 | 266.44 | 314.12 | 365. 75 | 421.27 |
| 40 | 223.36 | 267.20 | 314.95 | 366.64 | 422.23 |
| 41 | 224.06 | 267.96 | 315.78 | 367.53 | 423.19 |
| 42 | 224.76 | 268. 73 | 316.61 | 368.42 | 424.15 |
| 43 | 225.46 | 269.49 | 317.44 | 369.31 | 425.11 |
| 44 | 226.16 | 270.26 | 318.27 | 370.21 | 426.07 |
| 45 | 226.86 | 271.02 | 319.10 | 371.11 | 427.04 |
| 46 | 227.57 | 271.79 | 319.94 | 372.01 | 428.01 |
| 47 | 228.27 | 272.56 | 320.78 | 372.91 | 428.9 |
| 48 | 228.98 | 273.34 | 321.62 | 373.82 | 429.93 |


| lour angle <br> t | $10^{\mathrm{m}}$ | $11^{\mathrm{m}}$ | $12^{\mathrm{m}}$ | $13^{\mathrm{m}}$ | $14^{\mathrm{m}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 49 | 229.68 | 274.11 | 322.45 | 374.72 | 430.90 |
| 50 | 230.39 | 274.88 | 323.29 | 375.62 | 431.87 |
| 51 | 231.10 | 275.65 | 324.13 | 376.52 | 432.84 |
| 52 | 231.81 | 276.43 | 324.97 | 377.43 | 433.82 |
| 53 | 232.52 | 277.20 | 325.81 | 378.34 | 434.79 |
| 54 | 233.24 | 277.98 | 326.66 | 379.26 | 435.76 |
| 55 | 233.95 | 278.76 | 327.50 | 380.17 | 436.73 |
| 56 | 234.67 | 279.55 | 328.35 | 381.08 | 437.71 |
| 57 | 235.38 | 280.33 | 329.19 | 381.99 | 438.69 |
| 58 | 236.10 | 281.12 | 330.04 | 382.90 | 439.67 |
| 59 | 236.82 | 281.90 | 330.89 | 383.82 | 440.65 |
| 60 | 237.54 | 282.68 | 331.74 | 384.74 | 441.63 |


| $\underset{1^{\text {s }}}{\text { iff. }}$ | 0.1 | 0. 2 | 0.3 | 0. ${ }^{5}$ | 0.5 | 0.6 | 0.7 | 0. 8 | 0. 9 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0. 64 | 0.06 | 0.13 | 0.119 | 0.26 | 0.32 | 0. 38 | 0.45 | 0:" 51 | 0." 58 |
| 0.66 | 0.07 | 0.13 | 0.20 | 0.26 | 0.33 | 0.40 | 0.46 | 0.53 | 0.59 |
| 0.68 | 0.07 | 0.14 | 0.20 | 0.27 | 0.34 | 0.41 | 0.48 | 0.54 | 0.61 |
| 0.70 | 0.07 | 0.14 | 0.21 | 0.28 | 0.35 | 0.42 | 0.49 | 0.56 | 0.63 |
| 0. 72 | 0.07 | 0.14 | 0.22 | 0.29 | 0.36 | 0.43 | 0.50 | 0.58 | 0.65 |
| 0.74 | 0.07 | 0.15 | 0.22 | 0.30 | 0.37 | 0.44 | 0.52 | 0.59 | 0.67 |
| 0.76 | 0.08 | 0.15 | 0.23 | 0.30 | 0.38 | 0.46 | 0.53 | 0.61 | 0.68 |
| 0.78 | 0.08 | 0.16 | 0.23 | 0.31 | 0.39 | 0.47 | 0.55 | 0.62 | 0. 70 |
| 0. 80 | 0.08 | 0.16 | 0.24 | 0.32 | 0.40 | 0.48 | 0.56 | 0. 64 | 0.72 |
| 0.82 | 0.08 | 0.16 | 0.25 | 0.33 | 0.41 | 0.49 | 0.57 | 0.66 | 0. 74 |
| 0.84 | 0.08 | 0.17 | 0.25 | 0.34 | 0.42 | 0.50 | 0.59 | 0.67 | 0.76 |
| 0. 86 | 0.09 | 0.17 | 0.26 | 0.34 | 0.43 | 0.52 | 0.60 | 0.69 | 0.77 |
| 0.88 | 0.09 | 0.18 | 0.26 | 0.35 | 0.44 | 0.53 | 0.62 | 0. 70 | 0.79 |
| 0.90 | 0.09 | 0.18 | 0.27 | 0.36 | 0.45 | 0.54 | 0.63 | 0. 72 | 0. 81 |
| 0.92 | 0.09 | 0. 18 | 0.28 | 0.37 | 0.46 | 0.55 | 0.64 | 0. 74 | 0. 83 |
| 0.94 | 0.09 | 0.19 | 0.28 | 0.38 | 0.47 | 0.56 | 0.66 | 0. 75 | 0.85 |
| 0.96 | 0.10 | 0.19 | 0.29 | 0.38 | 0.48 | 0.58 | 0.67 | 0. 77 | 0.86 |
| 0.98 | 0.10 | 0.20 | 0.29 | 0.39 | 0.49 | 0.59 | 0.69 | 0.78 | 0.88 |


| Hour angle t | $15^{\text {m }}$ | $16^{\text {m }}$ | $17^{m}$ | $18^{m}$ | $19^{\text {m }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $0^{5}$ | 441 '. 63 | 502'. 46 | 567.' 19 | 635.' 85 | 708."42 |
| 1 | 442.62 | 503.50 | 568.30 | 637.02 | 709.6 |
| 2 | 443.60 | 504.55 | 569.42 | 638.20 | 710.9 |
| 3 | 444.58 | 505.60 | 570.53 | 639.38 | 712.15 |
| 4 | 445.56 | 506.65 | 571.65 | 640.56 | 713.3 |
| 5 | 446.55 | 507.70 | 572.76 | 641.74 | 714.64 |
| 6 | 447.54 | 508.76 | 573.88 | 642.93 | 715.89 |
| 7 | 448.53 | 509.81 | 575.00 | 644.11 | 717.14 |
| 8 | 449.51 | 510.86 | 576.12 | 645.30 | 718.3 |
| 9 | 450.50 | 511.92 | 577.24 | 646.48 | 719.6 |
| 10 | 451.50 | 512.38 | 578.36 | 647.67 | 720.8 |
| 11 | 452.49 | 514.93 | 579.48 | 648.86 | 722.15 |
| 12 | 453.48 | 515.09 | 580.61 | 650.05 | 723.40 |
| 13 | 454.48 | 516.15 | 581.73 | 651.24 | 724.6 |
| 14 | 455.47 | 517.21 | 582.86 | 652.43 | 725.9 |
| 15 | 456.47 | 518.27 | 583.99 | 653.62 | 727.1 |
| 16 | 457.47 | 519.34 | 585.12 | 654.80 | 728.43 |
| 17 | 458.47 | 520.40 | 586.24 | 656.01 | 729.6 |
| 18 | 459.47 | 521.47 | 587.38 | 657.21 | 730.9 |
| 19 | 460.47 | 522.53 | 588.51 | 658.40 | 732.2 |
| 20 | 461.47 | 523.60 | 589.64 | 659.60 | 733.4 |
| 21 | 462.48 | 524.67 | 590.77 | 660.80 | 734.7 |
| 22 | 463.48 | 525.74 | 591.91 | 662.00 | 736.0 |
| 23 | 464.48 | 526.81 | 593.05 | 663.20 | 737.2 |
| 24 | 465.49 | 527.89 | 594.18 | 664.40 | 738.5 |
| 25 | 466.50 | 528.96 | 595.32 | 665.61 | 739.8 |
| 26 | 467.51 | 530.03 | 596.46 | 666. 81 | 741.0 |
| 27 | 468.52 | 531.11 | 597.60 | 668.02 | 742.35 |
| 28 | 469.53 | 532.18 | 598. 74 | 669.22 | 743.62 |
| 29 | 470.54 | 533.26 | 599.88 | 670.43 | 744.8 |
| 30 | 471.55 | 534.33 | 601.03 | 671.64 | 746.1 |
| 31 | 472.57 | 535.41 | 602.17 | 672.85 | 747.4 |
| 32 | 473.58 | 536.50 | 603.32 | 674.06 | 748.7 |
| 33 | 474.60 | 537.58 | 604.47 | 675.28 | 750.0 |
| 34 | 475.62 | 538.67 | 605.61 | 676.49 | 751.2 |
| 35 | 476.64 | 539.75 | 606.76 | 677.70 | 752.5 |
| 36 | 477.65 | 540.83 | 607.91 | 678.92 | 753.8 |
| 37 | 478.67 | 541.91 | 609.06 | 680.13 | 755.12 |
| 38 | 479.70 | 543.00 | 610.22 | 681.35 | 756.4 |
| 39 | 480.72 | 544.09 | 611.37 | 682.57 | 757.6 |
| 40 | 481.74 | 545.18 | 612.53 | 683.79 | 758.9 |
| 41 | 482.77 | 546.27 | 613.68 | 685.01 | 760.2 |
| 42 | 483.79 | 547.36 | 614.84 | 686.23 | 761.5 |
| 43 | 484.82 | 548.45 | 616.00 | 687.46 | 762.8 |
| 44 | 485.85 | 549.55 | 617.15 | 688.68 | 764.12 |
| 45 | 486.88 | 550.64 | 618.31 | 689.91 | 765.4 |
| 46 | 487.91 | 551.73 | 619.48 | 691.13 | 766.7 |
| 47 | 488. 94 | 552.83 | 620.64 | 692.36 | 768.0 |
| 48 | 489.97 | 553.93 | 621.80 | 693.59 | 769.2 |


| Iour angle <br> $\mathbf{t}$ | $15^{\mathbf{m}}$ | $16^{\mathbf{m}}$ | $17^{\mathbf{m}}$ | $18^{m}$ | $19^{m}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 49 | 491.01 | 555.03 | 622.97 | 694.82 | 770.58 |
| 50 | 492.05 | 556.13 | 624.13 | 696.05 | 771.88 |
| 51 | 493.08 | 557.24 | 625.30 | 697.28 | 773.18 |
| 52 | 494.12 | 558.34 | 626.47 | 698.51 | 774.48 |
| 53 | 495.15 | 559.44 | 627.63 | 699.75 | 775.78 |
| 54 | 496.19 | 560.55 | 628.80 | 700.98 | 777.08 |
| 55 | 497.23 | 561.65 | 629.98 | 702.22 | 778.38 |
| 56 | 498.28 | 562.76 | 631.15 | 703.46 | 779.68 |
| 57 | 499.32 | 563.87 | 632.32 | 704.69 | 780.98 |
| 58 | 500.37 | 564.98 | 633.49 | 705.93 | 782.29 |
| 59 | 501.41 | 566.08 | 634.67 | 707.17 | 783.59 |
| 60 | 502.46 | 567.19 | 635.85 | 708.42 | 784.90 |


| Diff. per <br> $1^{s}$ | $0^{s .} 1$ | $0^{s .2}$ | $0^{s .3}$ | $0^{s .4}$ | $0^{s .5}$ | 0.5 | 0.7 | 0.8 | 0.9 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.98 | 0.410 | 0.20 | 0.29 | 0.39 | 0.49 | 0.59 | 0.69 | 0.78 | 0.88 |
| 1.00 | 0.10 | 0.20 | 0.30 | 0.40 | 0.50 | 0.60 | 0.70 | 0.80 | 0.90 |
| 1.02 | 0.10 | 0.20 | 0.31 | 0.41 | 0.51 | 0.61 | 0.71 | 0.82 | 0.92 |
| 1.04 | 0.10 | 0.21 | 0.31 | 0.42 | 0.52 | 0.62 | 0.73 | 0.83 | 0.94 |
| 1.06 | 0.11 | 0.21 | 0.32 | 0.42 | 0.53 | 0.64 | 0.74 | 0.85 | 0.95 |
| 1.08 | 0.11 | 0.22 | 0.32 | 0.43 | 0.54 | 0.65 | 0.76 | 0.86 | 0.97 |
| 1.10 | 0.11 | 0.22 | 0.33 | 0.44 | 0.55 | 0.66 | 0.77 | 0.88 | 0.99 |
| 1.12 | 0.11 | 0.22 | 0.34 | 0.45 | 0.56 | 0.67 | 0.78 | 0.90 | 1.01 |
| 1.14 | 0.11 | 0.23 | 0.34 | 0.46 | 0.57 | 0.68 | 0.80 | 0.91 | 1.03 |
| 1.16 | 0.12 | 0.23 | 0.35 | 0.46 | 0.58 | 0.70 | 0.81 | 0.93 | 1.04 |
| 1.18 | 0.12 | 0.24 | 0.35 | 0.47 | 0.59 | 0.71 | 0.83 | 0.94 | 1.06 |
| 1.20 | 0.12 | 0.24 | 0.36 | 0.48 | 0.60 | 0.72 | 0.84 | 0.96 | 1.08 |
| 1.22 | 0.12 | 0.24 | 0.37 | 0.49 | 0.61 | 0.73 | 0.85 | 0.98 | 1.10 |
| 1.24 | 0.12 | 0.25 | 0.37 | 0.50 | 0.62 | 0.74 | 0.87 | 0.99 | 1.12 |
| 1.26 | 0.13 | 0.25 | 0.38 | 0.50 | 0.63 | 0.76 | 0.88 | 1.01 | 1.13 |
| 1.28 | 0.13 | 0.26 | 0.38 | 0.51 | 0.64 | 0.77 | 0.90 | 1.02 | 1.15 |
| 1.30 | 0.13 | 0.26 | 0.39 | 0.52 | 0.65 | 0.78 | 0.91 | 1.04 | 1.17 |
| 1.32 | 0.13 | 0.26 | 0.40 | 0.53 | 0.66 | 0.79 | 0.92 | 1.06 | 1.19 |
| 1.34 | 0.13 | 0.27 | 0.40 | 0.54 | 0.67 | 0.80 | 0.94 | 1.07 | 1.21 |

But since $\varphi-\varphi_{c}=\left(\varphi-\varphi_{m}\right)+\left(\varphi_{m}-\varphi_{c}\right)$, we have :

$$
\begin{equation*}
\Delta \varphi_{c}=\Delta \varphi_{m}+\Delta \varphi \tag{8}
\end{equation*}
$$

where $\Delta \varphi_{m}$ stands for the constant error in the assumed value $\varphi_{m}$ and $\Delta \varphi$ for the difference (assumed minus computed) between the assumed value $\varphi_{m}$ and the computed value $\varphi_{c}$.

Accordingly the relations (6) and (7) can be re-written as :
$\Delta \varphi_{m}=-\Delta \varphi_{\mathrm{N}}+\sec a_{\mathrm{N}} \cdot \Delta h-\tan a_{\mathrm{N}} \cdot \cos \varphi_{m} \cdot \Delta t$
for north stars
and
$\Delta \varphi_{m}=-\Delta \varphi_{\mathrm{S}}-\sec a_{\mathrm{s}} \cdot \Delta h+\tan a_{\mathrm{S}} \cdot \cos \varphi_{m} \cdot \Delta t$
for south stars
Now eliminating $\Delta h$ between relations (9) and (10) and dividing both sides by $\cos a_{\mathrm{N}}+\cos a_{\mathrm{S}}$, we have, after simplification,
$\Delta \varphi_{m}=-\frac{\Delta \varphi_{\mathrm{N}} \cdot \cos a_{\mathrm{N}}+\Delta \varphi_{\mathrm{s}} \cdot \cos a_{\mathrm{S}}}{\cos a_{\mathrm{N}}+\cos a_{\mathrm{S}}}-15 \cdot \tan \frac{a_{\mathrm{N}}-a_{\mathrm{S}}}{2} \cdot \cos \varphi_{m} \cdot \Delta t$
where $\Delta t$ is in time.
Since $a_{\mathrm{N}}$ and $a_{\mathrm{S}}$ are small and do not differ considerably from each other, the first term of the R.H.S. of relation (11) which is the weighted mean of $-\Delta \varphi_{\mathrm{N}}$ and $-\Delta \varphi_{\mathrm{S}}$ (the respective weights being $\cos a_{\mathrm{N}}$ and $\cos a_{\mathrm{S}}$ ) can be easily reduced to the form :

$$
-\frac{\Delta \varphi_{\mathrm{N}}+\Delta \varphi_{\mathrm{S}}}{2}
$$

which is obviously the arithmetic mean of $-\Delta \varphi_{\mathrm{N}}$ and $-\Delta \varphi_{\mathrm{S}}$, without introducing any appreciable error unless $\Delta \varphi_{N}$ and/or $\Delta \varphi_{S}$ are abnormally large quantities. Thus the relation (11) ultimately takes the following form :

$$
\begin{equation*}
\Delta \varphi_{m}=-\frac{\Delta \varphi_{\mathrm{N}}+\Delta \varphi_{\mathrm{S}}}{2}-15 \cdot \tan \frac{a_{\mathrm{N}}-a_{\mathrm{S}}}{2} \cdot \cos \varphi_{m} \cdot \Delta t \tag{12}
\end{equation*}
$$

Now $a_{N}$ and $a_{\mathrm{S}}$ being again small and of the same sign, $\tan \frac{a_{N}-a_{\mathrm{s}}}{2}$ is necessarily of smaller order of magnitude and therefore can be easily and quickly computed by making use of only 3 -figure trigonometrical tables, without introducing any appreciable error unless $\Delta t$ is abnormally large, in which case 4 -figure tables may be necessary.

Accuracy : From the differential formulae (9), (10) and (12), it becomes obvious that the accuracy of the determination depends on :
(1) Precision of observed altitudes of stars, which is again dependent not only on the precision of the vertical circle readings of the theodolite in use, but also on that of the vertical collimation as well as the constant of atmospheric refraction. In the case of precision theodolites, viz. Wild T4, Wild T3 and the Geodetic Tavistock, the vertical angles are always read up to $0.1^{\prime \prime}$, though the accuracy actually obtainable with these instruments is generally of the order of $\pm 1.0^{\prime \prime}$ which is quite large. However in the case of Wild T4 the order of accuracy can be brought much lower down, say to about $0.2^{\prime \prime}$ or $0.1^{\prime \prime}$, by taking repeated readings of each vertical
setting, the actual intersections in each case being made by means of the moving micrometer only; but in the case of other theodolites it is unlikely to be improved by more than 50 per cent in any circumstances. As regards errors $\Delta h$ due only to the vertical collimation and the constant of atmospheric refraction, as already stated these are bound to remain practically constant during observations, and as such, cannot affect the final results appreciably. Thus if, apart from the possible errors arising from the vertical collimation and the constant of atmospheric refraction, there is an error of $\pm 1.0^{\prime \prime}$ in the two vertical circle readings for a star-pair, the corresponding error in the mean value of the latitude obtained will be of the order of $0.7^{\prime \prime}$, which is highly satisfactory.
(2) Precision of observed hour angles of stars, which is again dependent on the precision of the chronometer time - the error ( $\Delta t$ or $\Delta T$ ) in the latter being due to that in the assumed value $\lambda_{m}$ of longitude. But as already indicated while describing the procedure of observation, the error in the chronometer time can, with proper care, be reduced to a very small order of magnitude, so that ultimately its effect :

$$
-15 \cdot \tan \frac{a_{N}-a_{s}}{2} \cdot \cos \varphi_{m} \cdot \Delta t
$$

on the final value of latitude becomes either negligible on making $a_{\mathrm{N}}-a_{\mathrm{s}}$ sufficiently small, or easily computable to the desired accuracy by making use of the provisional value of the chronometer error $\Delta T(=\Delta t)$, as determined from longitude (time) observations by neglecting the effect of $\Delta \varphi$.

## LONGITUDE

Selection of stars : A star-pair, of magnitudes varying generally from 2.0 to 6.0 depending on the order of accuracy needed, one situated to the east and the other to the west but both within about 15 minutes of their time of elongation, must be selected for observation from the precise

Table 2
Giving ranges $r$ of declinations of stars at elongation at different latitudes for $t_{e}>20^{\circ}$ (numerically), $a_{e}>50^{\circ}$ (numerically) and $h_{e}>30^{\circ}$.

| $\varphi$ | $r$ |  |  |  | $\varphi$ | $\boldsymbol{r}$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $0^{\circ}$ |  | $0^{\circ}$ |  |  | $45^{\circ}$ | Between | $4{ }^{\circ}$ | and | $57^{\circ}$ |
| 5 | Between | $5.5{ }^{\circ}$ | and | $10^{\circ}$ | 50 | " | 52 | " | 60.5 |
| 10 | " | 10.5 | " | 20.5 | 55 | " | 56.5 | " | 64 |
| 15 | " | 16 | " | 31 | 60 | " | 61.5 | " | 67.5 |
| 20 | " | 21 | " | 43 | 65 | " | 66.5 | " | 71 |
| 25 | " | 26.5 | " | 46 | 70 | " | 71 | " | 75 |
| 30 | " | 31.5 | " | 48.5 | 75 | " | 76 | " | 78.5 |
| 35 |  | 36.5 | " | 51 | 80 | " | 80.5 | " | 82.5 |
| 40 | " | 42 |  | 54 | 85 | " | 85.5 | " | 86 |

fourth fundamental star catalogue, FK4. In addition, the altitudes of the star-pair should preferably be above $30^{\circ}$ and as nearly equal to each other as conveniently possible. While selecting a star-pair as above, use can however be made of table 2 which gives the various ranges of declinations and hour angles of stars at elongation for different latitudes, subject to the following conditions :
(a) Azimuths numerically $>50^{\circ}$;
(b) Hour angles numerically $<20^{\circ}$;
(c) Altitudes $>30^{\circ}$.

But since, for actual location in the fundamental star catalogue, the positions of stars must be known in terms of declinations and right ascensions, the range of hour angles required in table 2 cannot be usefully employed without reducing it to that of right ascensions by subtracting the values of hour angles from those of local sidereal times of observations derived directly from the corresponding standard mean times of observations by referring to table 3.

Table 3
Giving values of $R$
Change per day : 4 minutes
Day : from 0 to 24 hours U.T.

| Date | R |  | Date | R |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | h | m |  | h | m |
| Jan. 6 | 7 | 00 | July 8 | 19 | 02 |
| 22 | 8 | 03 | 23 | 20 | 01 |
| Feb. 6 | 9 | 03 | Aug. 7 | 21 | 00 |
| 21 | 10 | 02 | 23 | 22 | 03 |
| Mar. 8 | 11 | 01 | Sept. 7 | 23 | 02 |
| 23 | 12 | 00 | 22 | 0 | 02 |
| April 8 | 13 | 03 | Oct. 7 | 1 | 01 |
| 23 | 14 | 02 | 22 | 2 | 00 |
| May 8 | 15 | 01 | Nov. 7 | 3 | 03 |
| 23 | 16 | 01 | 22 | 4 | 02 |
| June 8 | 17 | 04 | Dec. 7 | 5 | 01 |
| 23 | 18 | 03 | 22 | 6 | 00 |

Observational programme : It is always essential to have a carefullymade programme lasting for about an hour in the case of first-order determinations, about half an hour in the case of second-order determinations and about a quarter of an hour in the case of third-order determinations, showing for the selected star-pair names, aspects, altitudes (computed from the approximate relation : $h=h_{e}+\left|t-t_{e}\right| \cos \delta$, where $\sin h_{e}=$ $\sin \varphi \cdot \operatorname{cosec} \delta$ and $\cos t_{e}=\tan \varphi \cdot \cot \delta$ ), and times (L.S.T.) of passages corresponding to the various azimuth positions of the star-pair (reduced from the relation :

$$
a=a_{e}-\tan a_{e} \cdot 2 \sin ^{2} \frac{t_{e}-t}{2} \operatorname{cosec} 1^{\prime \prime}
$$

where $\sin a_{e}=-\cos \delta \cdot \sec \varphi$, or $\operatorname{simply} \sin a \fallingdotseq-\cos \delta \cdot \sec \varphi$ ).

Observational equipment : The astronomical theodolite Wild T4 equipped with an impersonal micrometer, or any other theodolite of similar precision provided with a Talcott level, a good chronometer with chronograph equipment and a wireless set are all that are generally required for the first-order determinations. The Geodetic Tavistock with a Hunter Shutter and Wild T3 theodolites are usually used in the secondorder determinations, though the results in respect of the former are slightly better because of the impersonal shutter eye-piece employed. In the case of the third-order determinations, a Wild T2 or a Tavistock theodolite should be good enough.

Observational procedure : Before commencing actual observations, the line of horizontal collimation of the theodolite requires to be set in the meridian, correct to a quarter degree or so, knowing only the approximate value of the chronometer error and the chronometer time of meridian transit of a known circumpolar star. When necessary an approximate value of the latitude can also be deduced from the same observational results by making use of the relations : $\varphi=h_{0} \mp\left(90^{\circ}-\delta\right)$ or $\varphi=\delta+90^{\circ}-h_{0}$. The routine of observation is the same as in the case of latitude except that the system of time-recording in this case is more rigorous than in the case of latitude. While observing with the astronomical theodolite Wild T4 for first-order determinations, the procedure followed is similar to that of time-determination by meridian transits of stars observed only on one face of the theodolite at different vertical settings, their vertical circle readings being noted every time with adequate care. As usual the eye-piece micrometer of the theodolite in this case has to be turned through $90^{\circ}$, and the vertical bubble readings for each intersection taken from the Talcott level instead of the striding level. The case with the Geodetic Tavistock theodolite is similar. For its fuller description, reference may however be made to one of my previous papers : "A Method of Determining Astronomical Latitude and Longitude by Observing only Time and Horizontal Angles between Pairs of Stars", published in the Proceedings of the National Institute of Sciences of India, Vol. 26, A, No. 2, 1960. In the case of Wild T3, Wild T2 or Tavistock theodolites, the intersections of stars are carried out with the help of the horizontal wire and the instants of intersections noted with a tappet or simply a stop-watch, following the usual ear and eye method. As regards the series of altitudes measured, these need not be exactly symmetrical with respect to the meridian, but the number of intersections for the west star should in any case be the same as for the east star.

Computation : Let $h$ be the altitude of a star of declination $\delta$, corresponding to the hour angle $t$ when the star is near elongation, and $h_{e}$ be the altitude, corresponding to the azimuth $a_{e}$ and the hour angle $t_{e}$, when the star is at elongation.

Then on referring to fig. 2, we have :
$\sin h=\sin \varphi \cdot \sin \delta+\cos \varphi \cdot \cos \delta \cdot \cos t$
$\sin h_{e}=\sin \varphi \cdot \operatorname{cosec} \delta$
$\cos t_{e}=\tan \varphi \cdot \cot \delta$
and
$\sin a_{e}=-\cos \delta \cdot \sec \varphi$


Fig. 2
Denoting $h-h_{e}$ by $x$ and $t-t_{e}$ by $y$, we have from relations (1), (13) and (14) :

$$
\frac{\cos \left(y+t_{e}\right)-\cos t_{e}}{\cos t_{e}}=\frac{\sin \left(x+h_{e}\right)-\sin h_{e}}{\cos \varphi \cdot \cos \delta}
$$

or

$$
\cos y-\sin y \cdot \tan t_{e}-1=\frac{\sin x \cdot \cos h_{e}+\cos x \cdot \sin h_{e}-\sin h_{e}}{\cos \varphi \cdot \cos \delta \cdot \tan \varphi \cdot \cot \delta}
$$

or

$$
1-\cos y+\sin y \cdot \tan t_{e}=\left(1-\cos x-\sin x \cdot \cot h_{e}\right) \cdot \sec ^{2} \delta
$$

or

$$
\sin ^{2} \frac{y}{2}+\sin \frac{y}{2} \cdot \cos \frac{y}{2} \cdot \tan t_{e}=\sin ^{2} \frac{x}{2} \cdot \sec ^{2} \delta-\sin \frac{x}{2}
$$

$$
\cdot \cos \frac{x}{2} \cdot \cot h_{e} \cdot \sec ^{2} \delta
$$

$$
\begin{array}{r}
\left(\frac{y}{2}-\frac{y^{3}}{48}+\frac{y^{5}}{3840}-\ldots\right)^{2}+\left(\frac{y}{2}-\frac{y^{3}}{48}+\frac{y^{5}}{3840}-\ldots\right)\left(1-\frac{y^{2}}{8}+\frac{y^{4}}{384}-\ldots\right) \\
=\left(\frac{x}{2}-\frac{x^{3}}{48}+\frac{x^{5}}{3840}-\ldots\right)^{2} \cdot \sec ^{2} \delta-\left(\frac{x}{2}-\frac{x^{3}}{48}+\frac{x^{5}}{3840}-\ldots\right) \\
\left(1-\frac{x^{2}}{8}+\frac{x^{4}}{384}-\ldots\right) \cdot \cot h_{e} \cdot \sec ^{2} \delta
\end{array}
$$

or retaining terms up to only the $\overline{\text { th }}$ th powers of $x$ and $y$ : $y+\frac{y^{2}}{2} \cdot \cot t_{e}-\frac{y^{3}}{6}-\frac{y^{4}}{24} \cdot \cot t_{e}+\frac{y^{5}}{120}=-x \cdot \sec \delta+\frac{x^{2}}{2} \cdot \cot t_{e} \cdot \sec ^{2} \delta+$

$$
\begin{equation*}
+\frac{x^{3}}{6} \cdot \sec \delta-\frac{x^{4}}{24} \cdot \cot t_{e} \cdot \sec ^{2} \delta-\frac{x^{5}}{120} \cdot \sec \delta \tag{16}
\end{equation*}
$$

Now by the method of successive approximations, we obtain the required solution of the equation (16) in the following simplified form :

$$
\begin{aligned}
Y \text { in time }=-\frac{X}{15} & -\frac{X^{3}}{90} \cdot \sin ^{2} \delta \cdot \sin ^{2} 1^{\prime \prime}-\frac{X^{4}}{120} \cdot \sin ^{2} \delta \cdot \cot t_{e} \cdot \sin ^{3} 1^{\prime \prime} \\
& -\frac{X^{5}}{1800} \sin ^{2} \delta\left(9-\cos ^{2} \delta\right) \cdot \sin ^{4} 1^{\prime \prime}
\end{aligned}
$$

neglecting terms containing higher powers of X , where :

$$
\begin{equation*}
\mathbf{X}=x \cdot \sec \delta \tag{17}
\end{equation*}
$$

Since the term containing the 5 th power of $X$ in relation (17) is smaller than 0.01 s for observations of stars within about 20 minutes of their time of elongation we have finally :

$$
\begin{equation*}
t \text { in time }=t_{e} \text { in time }-\frac{\mathrm{X}}{15}-\mathbf{P}-\mathrm{Q} \tag{18}
\end{equation*}
$$

where

$$
P=0.2612 \cdot\left(\frac{X}{1000}\right)^{3} \cdot \sin ^{2} \delta
$$

which can be directly read from chart 3 without recourse to actual computations.

$$
\mathrm{Q}=0.364\left(\frac{\mathrm{P}}{10}\right)\left(\frac{\mathrm{X}}{1000}\right) \cot t_{e}
$$

which can also be directly read from chart 4 without requiring actual computations. And

$$
t_{e}>20^{\circ} \text { numerically }
$$

Since the hour angle $t$ of a star is equivalent to the chronometer time $T$ (L.S.T.) of observation minus the right ascension $\alpha$ of the star, relation (18) ultimately reduces to :

$$
\mathbf{T}=t_{e}+\alpha-\frac{\mathbf{X}}{15}-\mathbf{P}-\mathbf{Q}
$$

and

$$
\Delta t=\Delta(\mathbf{T}-\alpha)=\Delta \mathbf{T}
$$

Obviously the computed value $T_{c}$ of $T$, as derived from relation (19) by making use of the assumed values of $\varphi$ and $T$, will be in error by an amount equal to $\mathbf{T}-\mathrm{T}_{c}=\Delta \mathrm{T}_{c}=\Delta \mathbf{t}_{c}$ due to :
(1) Error $\Delta T_{m}$ in the assumed value $T_{m}$ of $T$, where $T-T_{m}=\Delta T_{m}$, a constant independent of the star of observation on account of the assumed value $\lambda_{m}$ of longitude being in error by the same amount $\Delta T_{m}$.
(2) Error $\Delta h$ in the observed altitudes of stars on account of the assumed values of both the vertical collimation and the constant of atmospheric refraction. The error is thus practically constant during observations unless the individual altitudes of the stars of observation differ considerably from each other due to faulty selection of stars.
(3) Error $\Delta \varphi$ in the assumed value of the unknown quantity $\varphi$.

But the above errors, when expressed in the form of differential formulae, take the following form :
$\Delta \mathrm{T}_{c}=\operatorname{cosec} a_{\mathrm{E}} \cdot \sec \varphi \cdot \Delta h-\cot a_{\mathrm{E}} \cdot \sec \varphi \cdot \Delta \varphi$
for east stars
and
$\Delta \mathrm{T}_{c}=-\operatorname{cosec} a_{\mathrm{w}} \cdot \sec \varphi \cdot \Delta h+\cot a_{\mathrm{w}} \cdot \sec \varphi \cdot \Delta \varphi$
for west stars
where $a_{E}$ and $a_{W}$ denote azimuthal angles as approximately determined from the relations :
$\sin a_{\mathrm{E}}=-\cos \delta \cdot \sec \varphi$ and $\sin a_{\mathrm{w}}=-\cos \delta \cdot \sec \varphi$
and suffixes ${ }_{E}$ and $w$ refer to east and west stars respectively.
But since $\mathrm{T}-\mathrm{T}_{c}=\left(\mathrm{T}-\mathrm{T}_{m}\right)+\left(\mathbf{T}_{m}-\mathrm{T}_{c}\right)$ we have :

$$
\Delta \mathrm{T}_{c}=\Delta \mathrm{T}_{m}+\Delta \mathrm{T}
$$

where $\Delta \mathrm{T}_{m}$ is the constant error in the assumed value $\mathrm{T}_{m}$ of the chronometer time $T$, and $\Delta T$ is the difference (assumed minus computed) between the assumed value $\mathrm{T}_{m}$ and the computed value $\mathrm{T}_{r}$.

Accordingly the relations (20) and (21) can be re-written as :
$\Delta \mathrm{T}_{m}=-\Delta \mathrm{T}_{\mathrm{E}}+\operatorname{cosec} \alpha_{\mathrm{E}} \cdot \sec \varphi \cdot \Delta h-\cot a_{\mathrm{E}} \cdot \sec \varphi \cdot \Delta \varphi$
for east stars
and
$\Delta \mathrm{T}_{m}=-\Delta \mathrm{T}_{\mathrm{w}}-\operatorname{cosec} a_{\mathrm{w}} \cdot \sec \varphi \cdot \Delta h+\cot a_{\mathrm{w}} \cdot \sec \varphi \cdot \Delta \varphi$
for west stars
Now eliminating $\Delta h$ between relations (22) and (23) and dividing both sides by $\sin a_{\mathrm{E}}+a_{\mathrm{W}}$, we have after simplification :
$\Delta \mathrm{T}_{m}=-\frac{\Delta \mathrm{T}_{\mathrm{E}} \cdot \sin a_{\mathrm{E}}+\Delta \mathrm{T}_{\mathrm{W}} \cdot \sin a_{\mathrm{W}}}{\sin a_{\mathrm{E}}+\sin a_{\mathrm{w}}}+\frac{1}{15} \cdot \tan \frac{a_{\mathrm{E}}-a_{\mathrm{W}}}{2} \cdot \sec \varphi_{m} \cdot \Delta \varphi$
Obviously the first term of the R.H.S. of relation (24) is the weighted mean of $-\Delta \mathrm{T}_{\mathrm{E}}$ and $-\Delta \mathrm{T}_{\mathrm{W}}$, the respective weights being $\sin a_{F}$ and $\sin a_{\mathrm{W}}$, where $a_{\mathrm{E}}$ and $a_{\mathrm{w}}$ lie between $50^{\circ}$ and $90^{\circ}$ and are easily computable by making use of 3 -figure trigonometrical tables only, unless $\Delta T_{E}$ and/or $\Delta \mathrm{T}_{\mathrm{W}}$ are abnormally large quantities. Again $a_{\mathrm{E}}-a_{\mathrm{W}}$ being small, $\boldsymbol{\operatorname { t a n }} \frac{a_{E}-a_{W}}{2}$ is of a smaller order of magnitude, and as such is easily computable by making use of 3 -figure trigonometrical tables only, without introducing any appreciable error unless $\Delta \varphi$ is abnormally large in which case 4 -figure tables may be necessary.

For longitude difference, however, it is necessary to determine at two stations the difference of correct local time of some event which can be perceived at both. The event is usually provided by some wireless rhythmic time-signals giving the seconds beats of the emitting local clock. These seconds beats are compared at the receiving station against those of the local clock whose error and rate have been known from astronomical observations, to determine the required longitude difference in time, after having corrected for propagation of signal and for a small error in the reputed time of emission of signal which is regularly published in the Admiralty Notices to Mariners and in the Bulletin Horaire. A detailed description of the procedure of determining accurately this longitude difference in time in the case of the Geodetic Tavistock fitted with shutter eyepiece, has been given in one of my previous papers referred to on page 111.

Accuracy : From the differential formulae (22), (23) and (24), it becomes obvious that the accuracy of the determination depends on :
(1) Precision of observed altitudes of stars, which is again dependent not only on the precision of the vertical circle readings of the theodolite in use, but also on that of the vertical collimation as well as the constant of atmospheric refraction. In the case of precision theodolites, viz. Wild T4, Wild T3 and the Geodetic Tavistock, the vertical angles are always read up to $0.1^{\prime \prime}$, though the accuracy actually obtainable with these instruments is generally of the order of $\pm 1.0^{\prime \prime}$, which is quite large. However in the case of Wild T4, the order of accuracy can be brought much lower down, say to $0.2^{\prime \prime}$ or $0.1^{\prime \prime}$, by taking repeated readings of each vertical settings, the actual intersections in each case being automatically made by the movement of the moving micrometer only. In the case of other theodolites, it is unlikely to be improved in any circumstances by more than 50 per cent, except in the case of the Geodetic Tavistock with which the results obtained are slightly better due to automatic time-recording carried out with a Hunter Shutter. As regards errors $\Delta h$, due only to the vertical collimation and the constant of atmospheric refraction, as already stated these are bound to remain practically constant during observations, and as such, cannot affect the final results appreciably. Thus if, apart from the possible errors arising from the vertical collimation and the constant of atmospheric refraction, there is an error of $\pm 1.0^{\prime \prime}$ in the two vertical circle readings for a star-pair, the corresponding error in the mean value of the chronometer time or longitude in time at a place of moderately high latitude is of the order of 0.05 s , which is highly satisfactory.
(2) Precision of the assumed value of $\varphi$. As already indicated while describing the procedure of observation the error $\Delta \varphi$ when very large can, with proper care, be reduced to a very small quantity, so that ultimately its effect $: \frac{1}{15} \cdot \tan \frac{a_{\mathrm{E}}-a_{\mathrm{W}}}{2} \cdot \sec \varphi \cdot \Delta \varphi$ on the final value of chronometer time or longitude in time becomes either negligible on making $a_{E}-a_{W}$ sufficiently small, or easily computable to the desired accuracy by making use of the provisional value of latitude error $\Delta \varphi$ as determined from latitude observations by neglecting the effect of $\Delta t$.

## CONCLUSION

Apart from its being extremely rapid and simple in its application, the method has many other practical advantages. Among these, the more important are :
(1) It is free from the error of vertical collimation, thereby obviating the necessity of observing on both faces of the theodolite as in other methods. This saves time and simplifies computations without sacrificing accuracy.
(2) It needs a much lesser number of stars, e.g. only two pairs of stars for the entire observations - one for latitude and the other for time or
longitude - and thus provides ample scope for the proper selection of stars whose declinations and right ascensions may be of dependable accuracy, which is not always possible in other methods.
(3) Altitudes of each star-pair being above $30^{\circ}$ and lying close to each other, it is practically free from the error of constant of atmospheric refraction - a factor of vital importance in astro-determinations by altitudes of stars.
(4) Computations involved are so simple that these can be easily worked out with the help of a few simple tables and charts specially designed for the purpose, and trigonometrical tables not exceeding 5 figures, except in the case of first-order determinations, when 6 -figure tables are needed in computing the values of $\mathrm{X}, \boldsymbol{t}_{e}$ and $h_{e}$.

The present method of astrofix is thus likely to prove more useful than any other method, especially in the field when both latitude and longitude are wanted in various survey operations including Laplace stations and astro-geodetic deflections at a large number of stations required in geodetic triangulations.

Specimen examples have been added at the end of this article in order to make the method clearer.

## Acknowledgement

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LATITUDE CHART 2
$E$ in seconds of arc
Argument : Am \& $h$


## EXAMPLES

## 1. Computation of latitude observations.

(a) First-order determination :

Formulae :
(1) $h_{0}=\varphi_{m} \pm\left(90^{\circ}-\delta\right) \quad$ for upper (lower) transit of north stars. $\boldsymbol{h}_{0}=90^{\circ}-\varphi_{m}+\delta \quad$ for south stars.
(2) $\mathbf{A}=\cos \varphi_{m} \cdot \cos \delta \cdot \sec h_{0}$
(3) $\varphi_{c}=h \mp\left(90^{\circ}-\delta\right)+A \cdot m-D+E$ for north stars near upper (lower) transit. $\varphi_{c}=\delta+\left(90^{\circ}-h\right)-\mathbf{A} \cdot \boldsymbol{m}+\mathbf{D}-\mathbf{E} \quad$ for south stars.
(4) $a_{N}^{\prime \prime}=-15 \cdot t^{s} \cdot \cos \delta \cdot \sec h_{0} \quad$ for north stars. $180^{\circ}-a_{s}^{\prime \prime}=-15 \cdot t^{\mathrm{s}} \cdot \cos \delta \cdot \sec h_{0} \quad$ for south stars.
(5) $\Delta \varphi_{m}=-\frac{\Delta \varphi_{\mathrm{N}}+\Delta \varphi_{\mathrm{s}}}{2}-15 \cdot \tan \frac{a_{\mathrm{N}}-a_{\mathrm{s}}}{2} \cdot \cos \varphi_{m} \cdot \Delta t$

| 1. Star : N/S | A Draconis : N | 12 Ophiuchi : S |
| :---: | :---: | :---: |
| 2. Altitude : $h^{(1)}$ | 51 ${ }^{\circ}$ 21' $41^{\prime \prime} 2$ | 57* $7^{\circ} 3^{\prime}$ 41:'6 |
| 3. L.S.T. of observation : T | $16^{\mathrm{h}} 08^{\mathrm{m}} \quad 27 \mathrm{~s} 36$ | $16^{\mathrm{h}} 18^{\mathrm{m}} 59: 70$ |
| 4. R.A. : $\alpha$ from FK4 | $16 \quad 28 \quad 05.46$ | $\begin{array}{lll}16 & 34 & 29.50\end{array}$ |
| 5. H.A. : t | - 011938.10 | - $0 \quad 15 \quad 29.80$ |
| 6. Decl. : $\delta$ from FK4 | $68{ }^{\circ} 50^{\prime}$ 57:3 | - 2 ${ }^{\circ} 14^{\prime}$ 56\%'6 |
| 7. $\cos \delta$ | 0.3608 | 0.9992 |
| 8. $\varphi_{m}$ | $30 \quad 19 \quad 10$ | $\begin{array}{llll}30 & 19 & 10\end{array}$ |
| 9. $\cos \varphi_{m}$ | 0.8632 | 0.8632 |
| 10. $h_{0}$ from formula (1) | $51 \quad 2813$ | $57 \quad 25 \quad 53$ |
| 11. sec $h_{0}$ | 1.6053 | 1.8577 |
| 12. A from formula (2) | 0.5000 | 1.6023 |
| 13. $m$ from table 1 | $756 \%$ \% | $471 \% 45$ |
| 14. A $\cdot \mathrm{m}$ | 378.3 | 755.4 |
| 15. $A \cdot m$ in min. and sec. | 6' 18:3 | 12' 35"4 |
| 16. D from chart ${ }^{(2)}$ | 0.4 | 2.2 |
| 17. E from chart (3) | 0.0 | 0.0 |
| 18. $\varphi_{c}$ from formula (3) | $30^{\circ} 188^{\prime} \quad 56^{\prime \prime} 4$ | $30^{\circ} 18^{\prime} \quad 48.6$ |
| 19. $\varphi_{m}-\varphi_{c}: \Delta \varphi \ldots$ | +13.6 | + 21.4 |
| 20. Mean of - $\Delta \varphi_{s}$ and - $\Delta \varphi_{y}$ | -17.5 |  |
| 21. $a$ in min. from formula (4) | $170^{\prime}$ | 431' |
| 22. $a$ in degr. and min. | 250 | 711 |
| 23. $\left(a_{\mathrm{N}}-a_{\mathrm{s}}\right) / 2$ | 210 |  |
| 24. $\tan \left(a_{N}-a_{S}\right) / 2$ | - 0.038 |  |
| 25. $\Delta t^{(4)}$ | - $1: 00$ |  |
| 26. Correction for $\Delta t$ from formula (5) | - 0\%\% |  |
| 27. $\Delta \varphi_{m}$ from formula (5) | $18: 0$ |  |
| 28. $\varphi_{m}+\Delta \varphi_{m}: \varphi \ldots . . . . . .$. | $30^{\circ} 18^{\prime} \quad 52^{\prime \prime} 0$ |  |

(1) Corrected for vertical inclination, collimation and refraction.
(2) Enter with A.m reading, ascend vertically to the curve on the left-hand side and then move horizontally towards the right. Enter again with $h_{0}$ reading, move horizontally to the curve on the right-hand side, ascend vertically up to the bounding horizontal line and move downwards following the slanting line (or the interpolated slanting line and move from the common position thus attained, descend vertically downwards to obtain the figure for $D$ from the horizontal scale on the right-hand side.
(3) Proceed as in (2) above after substituting $E$ for $D$.
(4) Only provisional value as obtained from time observations without considering the effect of $\Delta \varphi$.
(b) Second-order determination:

Proceed as in (a) above but the permissible error in the last decimal place for each figure in the present case is about $\pm 5$.
(c) Third-order determination :

## Formulae :

(1) $h_{0}=\varphi_{m} \pm\left(90^{\circ}-\delta\right) \quad$ for upper (lower) transit of north stars. $h_{0}=90^{\circ}-\varphi_{m}+\delta \quad$ for north stars.
(2) $\mathbf{A}=\cos \varphi_{m} \cdot \cos \delta \cdot \sec \boldsymbol{h}_{\mathbf{0}}$.
(3) $\varphi=h \mp\left(90^{\circ}-\delta\right)+A \cdot m$ for north stars near upper (lower) transit.
$\varphi=\delta+\left(90^{\circ}-h\right)-A \cdot m \quad$ for south stars.

| 1. Star : N/S | A Draconis : N | 12 Ophiuchi : S |
| :---: | :---: | :---: |
| 2. Observed altitude : $h$ | $51^{\circ} 21^{\prime} 40^{\prime \prime}$ | $57^{\circ} 13^{\prime} 40^{\prime \prime}$ |
| 3. L.S.T. of observation : T | $16^{\text {h }} 088^{\mathrm{m}} 30^{\text {m }}$ | $16^{\text {b }} 19^{\text {m }} \quad 03^{8}$ |
| 4. R.A. : $\alpha$ from any star almanac | $16 \quad 2805$ | $\begin{array}{llll}16 & 34 & 29\end{array}$ |
| 5. H.A. : $t$ | -0 19 35 | -0 15 26 |
| 6. Decl. : $\delta$ from any star almanac | $68^{\circ} 50^{\prime} 57^{\prime \prime}$ | $-2^{\circ} 14^{\prime} 57^{\prime \prime}$ |
| 7. $\cos \delta$ | 0.361 | 0.999 |
| 8. $\varphi_{m}$ | $30^{\circ} 18^{\prime} 000^{\prime \prime}$ | $30^{\circ} 18^{\prime \prime} 00^{\prime \prime}$ |
| 9. $\cos \varphi_{m}$ | 0.863 | 0.863 |
| 10. $h_{0}$ from formula (1) | $51^{\circ} 27^{\prime}$ | $57^{\circ} 27^{\prime}$ |
| 11. sec $h_{0}$ | 1.605 | 1.859 |
| 12. A from formula (2) | 0.500 | 1.602 |
| 13. $m$ from table 1 | $753^{\prime \prime}$ | $468{ }^{\prime \prime}$ |
| 14. A P m | 377 | 750 |
| 15. $A \cdot m$ in min. and sec. | $6^{\prime} 17^{\prime \prime}$ | $12^{\prime} 30^{\prime \prime}$ |
| 16. $\varphi$ from formula (3) | $30^{\circ} 18^{\prime \prime} 54 \prime$ | $30^{\circ} 18^{\prime} 53^{\prime \prime}$ |
| 17. Mean $\varphi$. | $\begin{array}{llll}30 & 18 & 54\end{array}$ |  |

## 2. Computation of clock correction observations.

(a) First-order determination :

## Formulae :

(1) $\cos t_{e}=\cot \delta \cdot \tan \varphi_{m}, t_{e}+v e$ for west stars and -ve for east stars.
(2) $\sin h_{e}=\operatorname{cosec} \delta \cdot \sin \varphi_{m}$
(3) $T_{c}=t_{e}+\alpha-\frac{X}{15}-P-Q$, last three terms beings of opposite sign in case of east stars.
(4) $\sin a=-\cos \delta \cdot \sec \varphi_{m}$
(5) $\mathrm{T}_{\mathrm{m}}=-\frac{\Delta \mathrm{T}_{\mathrm{E}} \cdot \sin a_{\mathrm{E}}+\Delta \mathrm{T}_{\mathrm{W}} \cdot \sin a_{\mathrm{W}}}{\sin a_{\mathrm{E}}+\sin a_{\mathrm{W}}}+\frac{1}{15}$.

$$
\tan \frac{a_{\mathrm{E}}-a_{\mathrm{W}}}{2} \cdot \sec \varphi_{m} \cdot \Delta \varphi
$$

| 1. Star : W/E | 23 Canum. <br> 23 Canum. | Lугæ : E |
| :---: | :---: | :---: |
| 2. Altitude : $h^{(1)}$ | $\begin{array}{llll} & 48^{\circ} & 16^{\prime} & 20 \% \\ \end{array}$ | $58^{\circ} 49^{\prime} 15^{\prime} 0$ |
| 3. L.S.T. of observation : $\mathrm{T}_{\boldsymbol{m}}$ | $\begin{array}{llll}16^{\mathrm{n}} & 40^{\mathrm{m}} & 15: 97\end{array}$ | $16^{\text {b }}$ 48 $8^{\text {m }}$ 01:94 |
| 4. R.A. : $\alpha$ from FK4 | $\begin{array}{lll}13 & 18 & 42.85\end{array}$ | $\begin{array}{lll}19 & 15 & 09.13\end{array}$ |
| 5. H.A. : $t$ | $\begin{array}{lll}3 & 21 & 33.12\end{array}$ | $\begin{array}{lll}-2 & 27 & 07.19\end{array}$ |
| 6. Decl. : $\delta$ from FK4 | $40^{\circ} 200^{\prime} \quad 31{ }^{\prime \prime} 8$ | $38^{\circ} 104^{\prime} 09 \% 9$ |
| 7. $\cot \delta$ | 1.177402 | 1.276750 |
| 8. $\varphi_{m}$ | $30^{\circ} 19^{\prime} 10^{\prime \prime}$ | $30^{\circ} 19^{\prime} 10^{\prime \prime}$ |
| 9. $\tan \varphi_{r}$ | 0.584808 | 0.584808 |
| 10. $\cos t_{e}$ from formula (1) | 0.688554 | 0.746654 |
| 11. $t_{e}$ in arc | $46^{\circ} 29^{\prime} \quad 03{ }^{\prime \prime} 3$ | -41 $41^{\prime} 55^{\prime \prime} 0$ |
| 12. $t_{c}$ in time | $3^{\text {h }} 05^{\text {m }} 56 \mathrm{E} 22$ |  |
| 13. $\sin \varphi_{m}$ | 0.504821 | 0.504821 |
| 14. cosec $\delta$ | 1.544757 | 1.621755 |
| 15. $\sin h_{e}$ from formula (2) | 0.779826 | 0.818696 |
| 16. $h_{e}$ | $51^{\circ} 14^{\prime} 40 \prime 6$ | $54^{\circ} 57^{\prime \prime} 16 \% 0$ |
| 17. $h-h_{e}: x$ | $-2 \quad 58 \quad 20.3$ | $3 \quad 51 \quad 59.0$ |
| 18. $x$ in seconds | -10700\%3 | 13919\%0 |
| 19. $\mathrm{sec} \delta$ | 1.312005 | 1.270222 |
| 20. $x$ sec $\delta$ | -14038:8 | $17680 \% 2$ |
| 21. X in time (opp. sign for east stars) | - $0^{\text {b }} 15^{\mathrm{m}} 35^{\text {s }} 92$ | $-0^{\mathrm{h}} 19^{\mathrm{m}} 38^{\mathrm{s}} 68$ |
| 22. $P$ from chart (2) (sign same as line 21) | $-0.30$ | $-0.55$ |
| 23. $Q$ from chart (3) (sign same as line 21) | 0.00 | -0.04 |
| 24. $\mathrm{T}_{e}+\alpha$ | $\begin{array}{lll}16 & 24 & 39.07\end{array}$ | $\begin{array}{lll}16 & 28 & 21.46\end{array}$ |
| 25. $\mathrm{T}_{c}$ from formula (3) | $\begin{array}{lll}16 & 40 & 15.29\end{array}$ | $\begin{array}{lll}16 & 48 & 00.73\end{array}$ |
| 26. $\mathrm{T}_{m}-\mathrm{T}_{c}: \Delta T$ | 0.68 | 1.21 |
| 27. $\cos \delta$ | 0.762 | 0.787 |
| 28. sec $\varphi_{m}$ | 1.158 | 1.158 |
| 29. $\sin a$ from formula (4) | 0.882 | 0.911 |
| 30. Weighted mean of $-\Delta \mathrm{T}_{\mathrm{w}}$ and <br> - $\Delta T_{\mathrm{E}}$ from formula (5) <br> 31. $a$ | $62^{\circ} 00^{\prime}-0895$ | $65^{\circ} 47^{\prime}$ |
| 32. $\left(a_{\mathrm{E}}-a_{\mathrm{W}}\right) / 2$ | - 154 | $65^{\circ} 4$ |
| 33. $\tan \left(a_{\mathrm{E}}-a_{W}\right) / 2$ | -0.0332 |  |
| 34. $\Delta \varphi^{(4)}$ | -18! 0 |  |
| 35. Correction for $\Delta \varphi$ from formula <br> (5) | $-0!05$ |  |
| 36. Clock correction : $\Delta \mathrm{T}_{m}$ from formula (5) | - 1900 |  |

(1) Corrected for vertical inclination, collimation and refraction.
(2) Proceed as in (2) of example 1 after substituting $X$ for $A \cdot m, \delta$ for $h_{0}$ and $P$ for $D$.
(3) Enter with $X$ reading, ascend vertically to the slanting line (or the interpolated slanting line) representing $P$ on the left-hand side and move horizontally towards the right. Enter again with $t_{e}$ reading, move horizontally to the curve on the right-hand side, ascend vertically up to the bounding horizontal line and then move downwards following the slanting line (or the interpolated slanting line), and finally from the common position thus attained, descend vertically downwards to obtain the figure for $Q$ from the horizontal scale on the right-hand side.
(4) Only provisional value as obtained from latitude observations without considering the effect of $\Delta t$.
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(b) Second-order determination :

Proceed as in (a) above but the permissible error in the last decimal place for each figure in the present case is about $\pm 5$.
(c) Third-order determination:

Formulae:
(1) $\cos \boldsymbol{t}_{e}=\cot \delta \cdot \tan \varphi_{m}, \boldsymbol{t}_{\boldsymbol{e}}+\boldsymbol{v e}$ for west stars and -ve for east stars.
(2) $\sin h_{e}=\operatorname{cosec} \delta \sin \varphi_{m}$
(3) $\mathrm{T}_{c}=t_{e}+\alpha-\frac{\mathrm{X}}{15}$, sign of $\frac{\mathrm{X}}{15}$ being opposite in case of east stars.
(4) $\Delta T_{m}=-\frac{\Delta T_{E}+\Delta T_{W}}{2}$

| 1. Star : W/E | 23 Canum. | Lyræ : E |
| :---: | :---: | :---: |
|  | Venat : W | $58^{\circ} 49^{\prime} 15^{\prime \prime}$ |
| 2. Observed altitude : $h$ | $48^{\circ} 16^{\prime} 20^{\prime \prime}$ | $50^{\circ} 49^{\prime} 15^{\prime \prime}$ |
| 3. L.S.T. of observation : $\mathrm{T}_{m}$ | $16^{\mathrm{h}} 40^{\mathrm{m}} 19 \mathrm{la}$ | $16^{\text {b }} 48^{\mathrm{m}} 0499$ |
| 4. R.A. : $\alpha$ from any star almanac | $\begin{array}{lll}13 & 18 & 42.9\end{array}$ | $\begin{array}{lll}19 & 15 & 09.1\end{array}$ |
| 5. H.A. : $t$ | $3{ }^{3} \quad 21 \quad 36.1$ | -2 2704.2 |
| 6. Decl. : $\delta$ from any star almanac | $40^{\circ} 20^{\prime} 32^{\prime \prime}$ | $38^{\circ} 04^{\prime} 10^{\prime \prime}$ |
| 7. $\cot \delta$ | 1.17740 | 1.27675 |
| 8. $\varphi_{m}$ | $30^{\circ} 18^{\prime} 00^{\prime \prime}$ | $30^{\circ} 18^{\prime} 000^{\prime \prime}$ |
| 9. $\tan \varphi_{m}$ | 0.58435 | 0.58435 |
| 10. $\cos t_{e}$ from formula (1) | 0.68801 | 0.74607 |
| 11. $t_{e}$ in arc | $46^{\circ} 31^{\prime} 38^{\prime \prime}$ | - $41^{\circ} 44^{\prime} 56^{\prime \prime}$ |
| 12. $t_{e}$ in time | $3^{\text {h }} 06^{\mathrm{m}} 0685$ | - $2^{\mathrm{h}} \mathbf{4 6}^{\mathrm{m}} 59: 7$ |
| 13. $\sin \varphi_{m}$ | 0.50453 | 0.50453 |
| 14. cosec $\delta$ | 1.54476 | 1.62175 |
| 15. $\sin h_{e}$ from formula (2) | 0.77938 | 0.81822 |
| 16. $h_{e}$ | $51^{\circ} 12^{\prime} 14^{\prime \prime}$ | $54^{\circ} 54^{\prime} 25^{\prime \prime}$ |
| 17. $h-h_{e}: x$ | $255 \quad 54$ | 35450 |
| 18. $x$ in seconds | - 10554" | 14090" |
| 19. sec $\delta$ | 1.31201 | 1.27022 |
| 20. $x \sec \delta: \mathrm{X}$ | -13847' | 17897" |
| 21. $X$ in time (opp. sign for east stars) |  | - $0^{\mathrm{h}} 19^{\mathrm{m}} 53 \mathrm{~s}$. |
| 22. $\mathrm{T}_{c}$ from formula (3) | $16 \quad 40 \quad 12.5$ | $\begin{array}{llll}16 & 48 & 02.4\end{array}$ |
| 23. $\mathbf{T}_{m}-\mathrm{T}_{c}=\Delta \mathbf{T} \ldots \ldots \ldots \ldots \ldots$ | 6.5 | 2.4 |
| 24. Mean clock correction : $\Delta \mathrm{T}_{\mathrm{m}}$ from formula (4) | -4.4 |  |

## Conversion of time

L.S.T. = standard time plus (local long. minus standard long.) in time plus $R$ (interpolated for the date and the fraction of the day).
$N . B$. - The table is correct to 2 minutes for all years. For better accuracy, however, the tabular values may be taken to be true for :
(1) $1962,1966,1970,1974 \ldots$ i.e. two years after the leap years.
(2) All leap years after increasing the values by 2 minutes except for the months of January and February in which case the tabular values are to be decreased by 2 minutes.
(3) Years following the leap years after increasing by 1 minute. and (4) Remaining years after decreasing by 1 minute.

